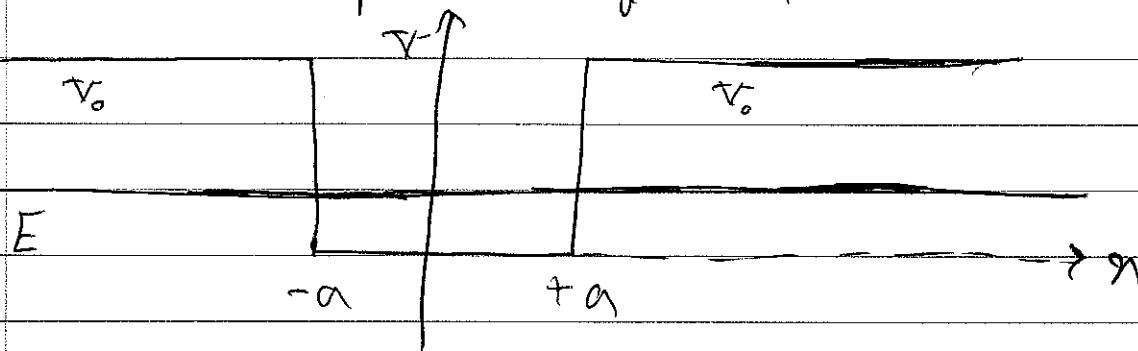


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Tunneling:

Consider the square well potential:



A state with energy  $E < V_0$  (there is always at least one such state) is a bound state and stable,

i.e.  $|\Psi(x,t)|^2$  is independent of time. Note that

$\Psi$  is an exponentially damped function at  $x > +a$ ,

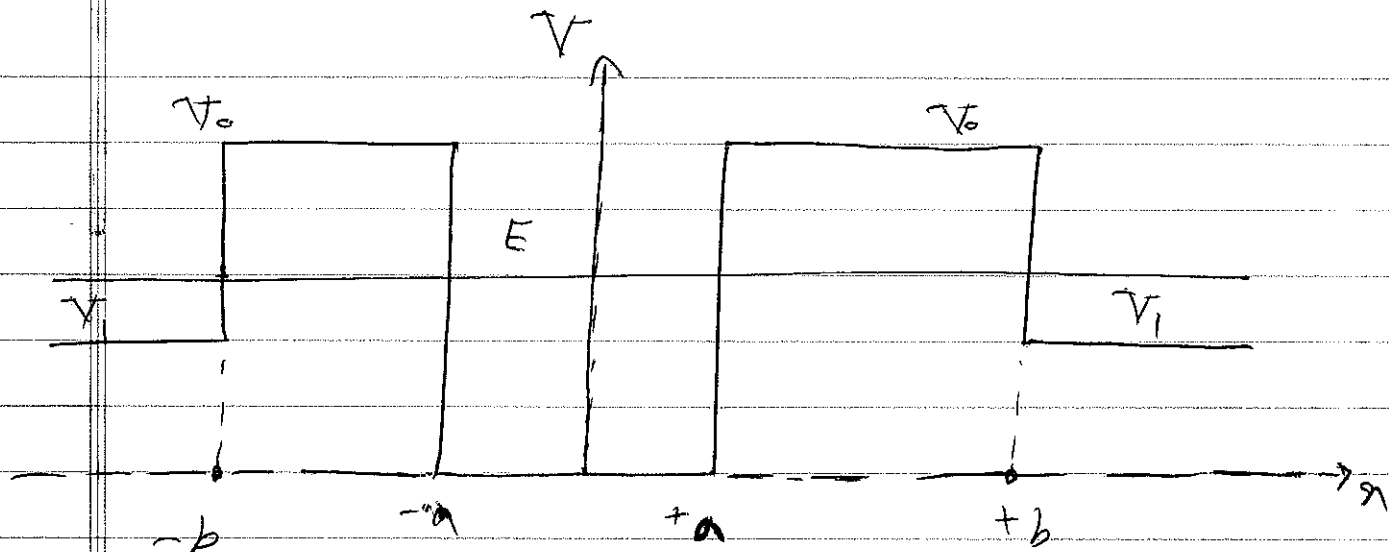
$x < -a$ .

Now, let's assume that the potential is lowered from

$V_0$  to  $V_1 < V_0$  for  $x > +b$ ,  $x < -b$  where  $|b| > |a|$ .

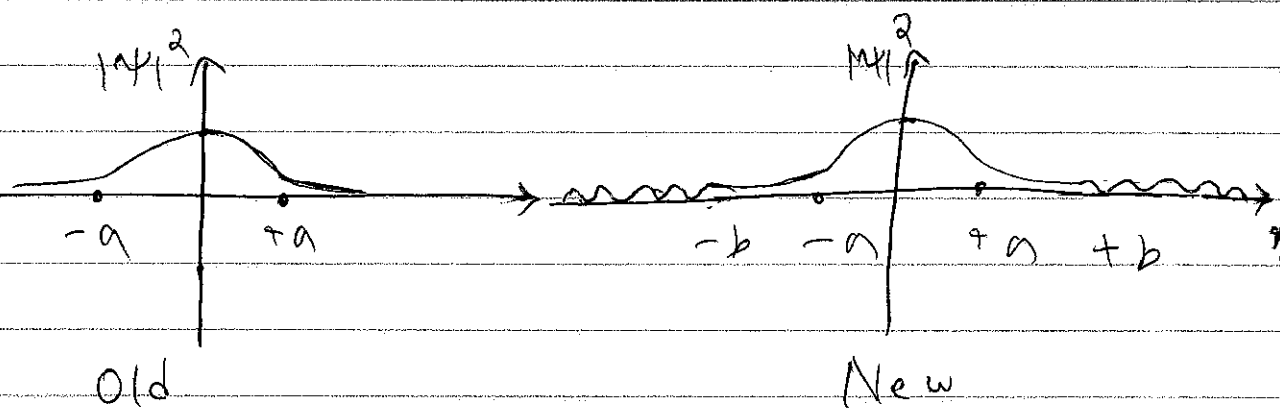
The question is what will happen to the state

$\Psi$  that was a stable eigenstate of the Hamiltonian.



The Hamiltonian has an eigenstate with energy  $E$  in this case too. In fact, the energy spectrum is continuous for  $E > V_1$ .

However,  $\psi$  cannot be the eigenstate with energy  $E$  now. For the new potential, the energy eigenstate is oscillatory at  $x > +b$ ,  $x < -b$ , while for the old potential the eigenstate is exponentially damped.



Therefore, the eigenstate of the old potential will be a superposition of the eigenstates of the new potential. It will be time-dependent and not stable any more. It is a metastable state that leaks outside the  $(-a, +a)$  interval.

This is total against classical intuition. In classical mechanics, if  $E < V$ , the particle cannot climb over the barrier and will always remain within the  $(-a, +a)$  interval.

Why doesn't this work in quantum mechanics?

It all goes down to nonzero commutator of  $X$  and  $H$ . In a given energy eigenstate,  $\Delta X \neq 0$  because the energy eigenstate is not an  $X$  eigenstate.

Hence there is an uncertainty in the position and the particle has a nonzero probability to

to be beyond the well.

On the other hand, if the particle is in an eigenstate of  $X$ , it will not have definite energy. In this case  $\Delta H \neq 0$ , and the components with  $E > V$ , can take the particle over the barrier.

Therefore a particle localized inside the  $[a, +a]$  interval, is in an unstable state. If one waits for a sufficiently long time, there is a probability to find the particle at  $x > +b$  or  $x < -b$ . This is called tunneling, which is a purely quantum mechanical effect.

Lets consider two important examples of tunneling.

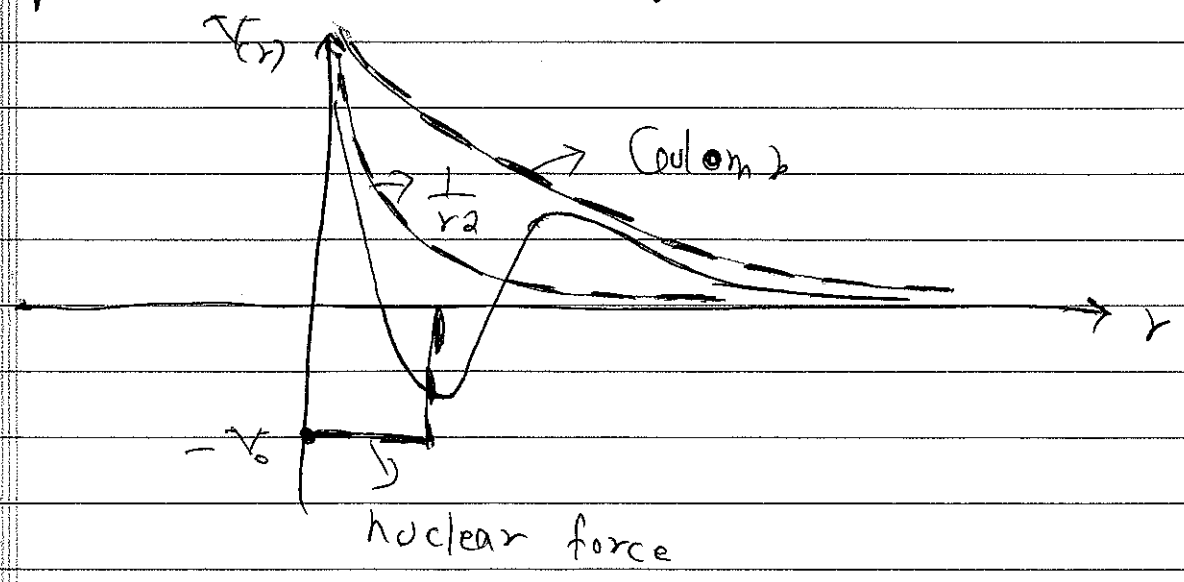
### Nuclear $\alpha$ Decay:

$\alpha$  particle consists of two protons and two neutrons.

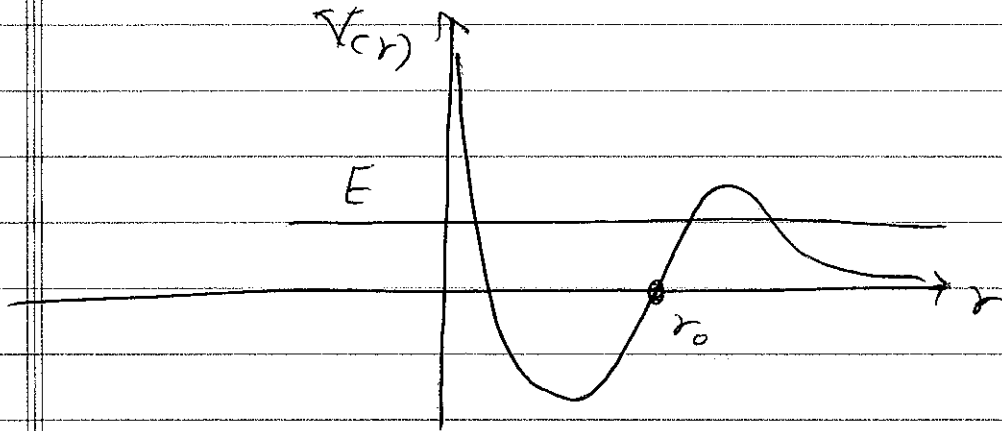
Inside a large nucleus,  $\alpha$  particle is subject to various potentials:

- 1. Attractive short range potential from nuclear forces,
- 2. Repulsive Coulomb potential from other protons in the nucleus.
- 3. Repulsive  $\frac{1}{r^2}$  potential for states with non zero angular momentum.

The combined effect leads to the following potential for the  $\alpha$  particle:



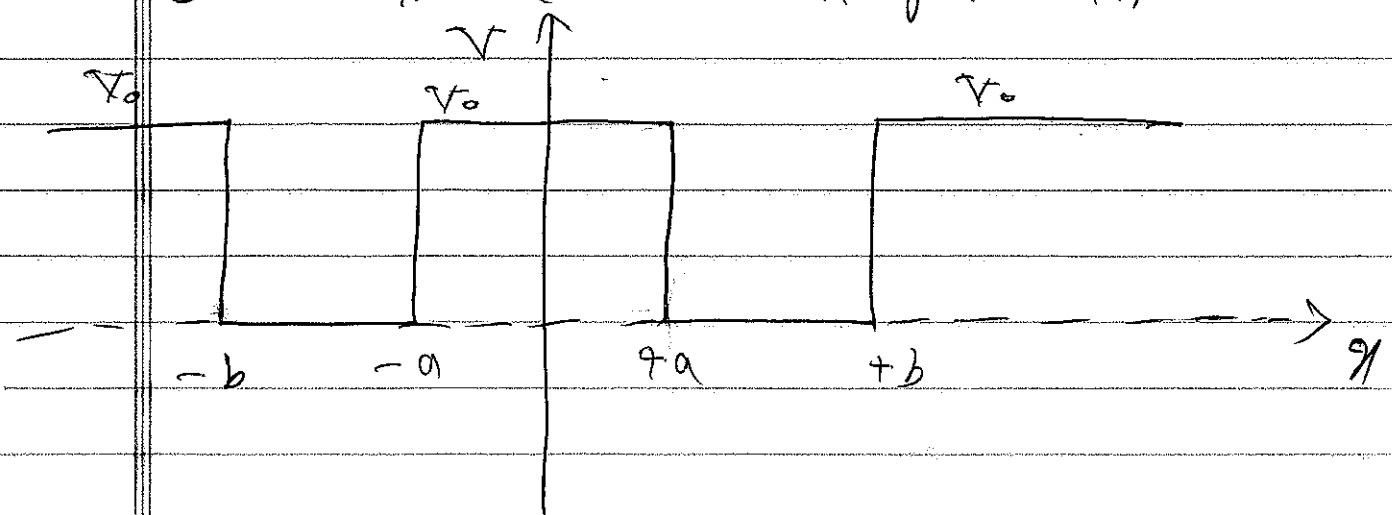
Note that for a positive energy, the eigenstate has oscillatory behavior at large  $r$ :



Therefore, an  $\alpha$  particle localized at  $r < r_0$  (contained in the nucleus) can leak out. This will result in a decay of the nucleus.

### Tunneling and Removal of Degeneracy:

Consider the double well potential;



Classically, the particle exists in only one

of the intervals  $[-b, -a]$  or  $[+a, +b]$ . We therefore have a two-fold degeneracy.

However, in Quantum mechanics tunneling effect removes this degeneracy. The particle can be found in both wells, and the probability find it in  $[-b, -a]$  interval is the same as the probability to find it in  $[+a, +b]$  interval.

Instead of the classical degeneracy, we have two states with different energies: one is an even function of  $x$  (lower energy state) and one is an odd function of  $x$  (higher energy state).